



FIG. 2. The differences of the principal specific susceptibilities of a zinc crystal as a function of the applied magnetic field strength at  $T = 4.2^\circ \text{K}$ . At  $\theta = 20^\circ$ : (a)  $P = 0 \text{ kg/cm}^2$ ; (b)  $P \sim 1500 \text{ kg/cm}^2$ ; (c) pressure removed; (d) pressure  $P \sim 1500 \text{ kg/cm}^2$  reapplied; (e) reapplied pressure removed. At  $\theta = 80^\circ$ : (a)  $P = 0 \text{ kg/cm}^2$ ; (b)  $P \sim 1500 \text{ kg/cm}^2$ ; (c) pressure removed.

Using the data obtained, one can estimate the number of electrons in the group responsible for the investigated longest-period component of the de Haas-van Alphen effect in Zn. If  $n$  is the number of electrons in the given group, and  $\Omega$  is the volume bounded by the corresponding Fermi surface in momentum space, then  $n = \Omega V/h^3$ , where  $V$  is the volume of the metal and  $h$  is Planck's constant. The equation  $\Omega = \alpha S^{3/2}$  gives the relation between the volume  $\Omega$  and the extremal cross sectional area of the Fermi surface, and  $S = A/\Delta$  [where  $\Delta$  is the period of the  $\chi(1/H)$  curve]. Here  $\alpha$  is a form factor that depends in general on the orientation of the field vector relative to the crystal axis. From these formulas we obtain

$$\delta n/n = \delta V/V - 3/2 (\delta \Delta/\Delta) + \delta \alpha/\alpha.$$

In our case, in spite of the significant anisotropy of the compressibility of Zn, the fractional changes of the periods for  $\theta = 20^\circ$  and  $\theta = 80^\circ$  are close (0.52 and 0.43 respectively). We can thus assume that the shape of the Fermi surface changes but little, and set  $\delta \alpha = 0$ . Taking next  $(\delta \Delta/\Delta)_{av} = 0.47$  and  $(\delta V/V) \approx 3 \cdot 10^{-3}$ , we obtain  $\delta n/n \approx -0.7$ . In other words, the number of electrons responsible for the long-period component in the  $\chi(1/H)$  curve for Zn is about 70%. For comparison, we note that a similar evaluation undertaken for Bi [taking  $\delta \Delta/\Delta = 0.13$ , obtained from the  $\chi(1/H)$  curves for  $\theta = -70^\circ$ ] gives only a 10% decrease of the number of electrons responsible for the de Haas-van Alphen effect at a pressure of about  $1500 \text{ kg/cm}^2$ . This estimate is much rougher than the one performed here for zinc, owing to the way the anisotropy of the Fermi surface of Bi changes under pressure.<sup>2</sup>

Finally, it is interesting to note that the work per atom done by the external forces in compressing the crystal is  $2 \times 10^{-6} \text{ ev}$ . To some extent this value characterizes the magnitude of the change of the binding forces in the Zn lattice under a pressure of about  $1500 \text{ kg/cm}^2$ . It is interesting that such a small change in the binding forces leads to so large a change in the parameters of the de Haas-van Alphen effect, owing to the anomalously small group of electrons in zinc with the Fermi energy  $E_0 \sim 10^{-2} \text{ ev}$ .

In conclusion the authors consider it their pleasant duty to thank I. M. Lifshitz for discussing the results of the work.

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<sup>2</sup> Verkin, Dmitrenko, and Lazarev, *J. Exptl. Theoret. Phys.* **31**, 538 (1956); *Soviet Phys. JETP* **4**, 432 (1957).

<sup>3</sup> B. I. Verkin, *Dokl. Akad. Nauk SSSR* **81**, 529 (1951); B. I. Verkin and I. F. Mikhailov, *J. Exptl. Theoret. Phys. (U.S.S.R.)* **24**, 342 (1953).

<sup>4</sup> S. Sidoriak and D. Robinson, *Phys. Rev.* **75**, 118 (1949); J. Marcus, *Phys. Rev.* **84**, 787 (1951); F. Donahoe and F. Nix, *Phys. Rev.* **95**, 1395 (1954); T. Berlincourt and M. Steele, *Phys. Rev.* **95**, 1421 (1954).

<sup>5</sup> B. I. Verkin and I. M. Dmitrenko, *Izv. Akad. Nauk SSSR, ser. fiz.* **19**, 409 (1955).

<sup>6</sup> B. I. Verkin, I. F. Mikhailov, *J. Exptl. Theoret. Phys. (U.S.S.R.)* **25**, 471 (1953).